

Torque on a magnetic dipole

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We saw that the force acting on a magnetic dipole is

$$\begin{aligned}\bar{F} &= -\nabla U \quad \text{with} \quad U = -\bar{m} \cdot \bar{B} \\ \bar{F} &= \nabla (\bar{m} \cdot \bar{B}) = \bar{m} \times (\nabla \times \bar{B}) + \bar{B} \times (\nabla \times \bar{m}) + (\bar{m} \cdot \nabla) \bar{B} \\ &\quad + (\bar{B} \cdot \nabla) \bar{m} \\ &= \bar{m} \times (\underbrace{\nabla \times \bar{B}}_{=0 \text{ if } \bar{j}=0}) + (\bar{m} \cdot \nabla) \bar{B}\end{aligned}$$

The force is then zero if \bar{B} is constant.

Here we want to determine the torque acting on the magnetic dipole. One can start from the vector potential due to a dipole placed at \bar{x}_0 .

$$\bar{A}(\bar{x}) = \frac{\mu_0}{4\pi} \frac{\bar{m} \times (\bar{x} - \bar{x}_0)}{|\bar{x} - \bar{x}_0|^3} \quad \longrightarrow \quad \bar{B} = \nabla \times \bar{A}$$

rem

$$\nabla \times (\bar{a} \times \bar{b}) = (\bar{b} \cdot \nabla) \bar{a} - (\bar{a} \cdot \nabla) \bar{b} + \bar{a} (\nabla \cdot \bar{b}) - \bar{b} (\nabla \cdot \bar{a})$$

$$\nabla \times \left(\bar{m} \times \frac{(\bar{x} - \bar{x}_0)}{|\bar{x} - \bar{x}_0|^3} \right) = -(\bar{m} \cdot \nabla) \left(\frac{\bar{x} - \bar{x}_0}{|\bar{x} - \bar{x}_0|^3} \right) + \bar{m} \nabla \cdot \left(\frac{\bar{x} - \bar{x}_0}{|\bar{x} - \bar{x}_0|^3} \right)$$

$$\begin{aligned}\bar{B}(\bar{x}) &= \frac{\mu_0}{4\pi} \left[\bar{m} \nabla \cdot \frac{\bar{x} - \bar{x}_0}{|\bar{x} - \bar{x}_0|^3} - (\bar{m} \cdot \nabla) \frac{\bar{x} - \bar{x}_0}{|\bar{x} - \bar{x}_0|^3} \right] \\ &= \frac{\mu_0}{4\pi} \left[-\bar{m} \nabla \cdot \left(\nabla \frac{1}{|\bar{x} - \bar{x}_0|} \right) - (\bar{m} \cdot \nabla) \frac{\bar{x} - \bar{x}_0}{|\bar{x} - \bar{x}_0|^3} \right]\end{aligned}$$

rem

$$\nabla (\bar{a} \cdot \bar{b}) = \bar{a} \cdot \nabla \bar{b} + \bar{b} \cdot \nabla \bar{a}$$

$$\bar{B}(\bar{x}) = \frac{\mu_0}{4\pi} \left[-\bar{m} \Delta \frac{1}{|\bar{x} - \bar{x}_0|} - \nabla \left(\frac{\bar{m} \cdot (\bar{x} - \bar{x}_0)}{|\bar{x} - \bar{x}_0|^3} \right) \right]$$

From Ampere's equation in magnetostatics one has

$$\nabla \times \bar{B} = \mu_0 \bar{J} \rightarrow \bar{J} = \frac{1}{\mu_0} \nabla \times \bar{B}$$

$$\bar{J}(\bar{x}) = \frac{1}{4\pi} \left[-\nabla \times \left(\bar{m} \underbrace{\Delta \frac{1}{|\bar{x} - \bar{x}_0|}}_{-4\pi \delta^{(3)}(\bar{x} - \bar{x}_0)} \right) - \underbrace{\nabla \times \nabla}_{=0} \frac{\bar{m} \cdot (\bar{x} - \bar{x}_0)}{|\bar{x} - \bar{x}_0|^3} \right]$$

$$= \nabla \times \left(\bar{m} \delta^{(3)}(\bar{x} - \bar{x}_0) \right)$$

rem $\nabla \times (f \bar{A}) = f \nabla \times \bar{A} - \bar{A} \times \nabla f$

$$\bar{J}(\bar{x}) = -\bar{m} \times \nabla \delta^{(3)}(\bar{x} - \bar{x}_0)$$

CURRENT
DENSITY OF
A MAGNETIC
DIPOLE

The torque that we want to evaluate is

$$\bar{\Gamma} = \int \bar{x} \times d\bar{F} = \int d^3x \bar{x} \times (\bar{J} \times \bar{B})$$

external \bar{B} ,
not the \bar{B} due
to \bar{J}

$$= - \int d^3x \bar{x} \times \left[(\bar{m} \times \nabla \delta^{(3)}(\bar{x} - \bar{x}_0)) \times \bar{B}(\bar{x}) \right]$$

$$= \int d^3x \bar{x} \times \left[\bar{B}(\bar{x}) \times (\bar{m} \times \nabla \delta^{(3)}(\bar{x} - \bar{x}_0)) \right]$$

rem

$$\bar{a} \times (\bar{b} \times \bar{c}) = (\bar{a} \cdot \bar{c}) \bar{b} - (\bar{a} \cdot \bar{b}) \bar{c}$$

$$\bar{\Gamma} = \int d^3x \bar{x} \times \left[(\bar{B} \cdot \nabla \delta^{(3)}(\bar{x} - \bar{x}_0)) \bar{m} - (\bar{B} \cdot \bar{m}) \nabla \delta^{(3)}(\bar{x} - \bar{x}_0) \right]$$

Consider now a component of the torque

$$\Gamma_i = \int d^3x \varepsilon_{ijk} x_j \left[B_\ell \partial_\ell \delta^{(3)}(\bar{x} - \bar{x}_0) m_k - m_\ell B_\ell \partial_k \delta^{(3)}(\bar{x} - \bar{x}_0) \right]$$

$$= - \int d^3x \varepsilon_{ijk} \partial_\ell (x_j B_\ell) \delta^{(3)}(\bar{x} - \bar{x}_0) m_k$$

$$+ \int d^3x \varepsilon_{ijk} \partial_k (x_j B_\ell) \delta^{(3)}(\bar{x} - \bar{x}_0) m_\ell$$

BY
PARTS

In the integration by parts, the surface terms vanish since we are integrating over \mathbb{R}^3 and the delta function is zero everywhere except for \vec{x}_0 .

$$\begin{aligned}
 \vec{\tau}_i &= \int d^3x \delta^{(3)}(\vec{x} - \vec{x}_0) \varepsilon_{ijk} \left[m_\ell \partial_k - m_k \partial_\ell \right] x_j B_\ell \\
 &= \int d^3x \delta^{(3)}(\vec{x} - \vec{x}_0) \varepsilon_{ijk} \left[m_\ell \delta_{jk} B_\ell + m_\ell x_j \partial_k B_\ell \right. \\
 &\quad \left. - m_k \delta_{\ell j} B_\ell - m_k x_j \underbrace{\partial_\ell B_\ell}_{\nabla \cdot \vec{B} = 0} \right] \\
 &= \int d^3x \delta^{(3)}(\vec{x} - \vec{x}_0) \varepsilon_{ijk} x_j \underbrace{\partial_k (\vec{m} \cdot \vec{B})}_{\partial_k (\vec{m} \cdot \vec{B})} - \int d^3x \delta^{(3)}(\vec{x} - \vec{x}_0) \varepsilon_{ijk} m_k B_j \\
 &= \left[\vec{m} \times \vec{B}(\vec{x}_0) \right]_i + \left[\vec{x}_0 \times \nabla (\vec{m} \cdot \vec{B}) \right]_i
 \end{aligned}$$

We saw that the force acting on a dipole is

$$\begin{aligned}
 \vec{F} &= -\nabla (-\vec{m} \cdot \vec{B}) = \nabla (\vec{m} \cdot \vec{B}) \\
 \vec{x}_0 \times \nabla (\vec{m} \cdot \vec{B}) &= \vec{x}_0 \times \vec{F}(\vec{x}_0)
 \end{aligned}$$

So that finally

$$\vec{\tau} = \vec{m} \times \vec{B}(\vec{x}_0) + \vec{x}_0 \times \vec{F}(\vec{x}_0)$$

Consequently if the dipole is located at the origin of the frame of reference one finds

$$\vec{x}_0 = 0 \longrightarrow \boxed{\vec{\tau} = \vec{m} \times \vec{B}} \quad \text{TORQUE ACTING ON A MAGNETIC DIPOLE}$$