Electrostatic potential

Tuesday, January 4, 2022

4:03 AM

Remember the electrostatics equations

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$
 CAUSS' EQUATION FOR \vec{E}

The second equation implies that electric field is the gradient of a scalar function, which is called scalar potential or electrostatic potential (or simply potential when there are no ambiguities)

$$\nabla \times E = 0$$
 $\longrightarrow E = -\nabla \varphi$ (the minus sign)

The potential is defined up to an additive constant

$$E = -\nabla \varphi = -\nabla (\varphi + \varphi_o)$$
 constant

One can then rewrite Gauss' law in terms of the potential, in this way one finds

Poisson's equation

$$\nabla \cdot \vec{E} = \nabla \cdot (-\nabla \varphi) = -\nabla \cdot \nabla \varphi = -\Delta \varphi = \frac{\rho}{\varepsilon_o}$$

$$\Delta \varphi = -\frac{\rho}{\varepsilon_o}$$
POISSON'S EQUATION

Formal solution of Poisson's equation

Poisson's equation has a useful solution for a charge distribution localized in a finite volume. In this solution the constant ambiguity in the potential is fixed by requiring that the potential vanishes at infinity far away from the source charges.

$$\varphi(\bar{r}) = \frac{1}{4\pi \epsilon_o} \int d^3r' \frac{p(\bar{r}')}{|\bar{r} - \bar{r}'|}$$

The equation above is indeed the solution of Poisson's equation since

$$\Delta_{r} \varphi(\bar{r}) = \frac{1}{4\pi \epsilon_{o}} \int_{\sigma}^{3r} \rho(\bar{r}') \Delta_{r} \frac{1}{|\bar{r} - \bar{r}'|}$$

$$-4\pi \delta^{(3)}(\bar{r} - \bar{r}')$$

$$= -\frac{1}{\epsilon_{o}} \int_{\sigma}^{3r} \rho(\bar{r}') \delta^{(3)}(\bar{r} - \bar{r}')$$

$$= -\frac{\rho(\bar{r})}{\epsilon_{o}}$$

$$Q.E.D.$$

This is formally all what we need to study electrostatic situations with charges that are confined in a finite volume. Unfortunately, the evaluation of the integral needed to determine the potential can in general be very complicated.