

# Charged particle in a magnetic field

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The goal of this section is to find a function  $L$  such that the corresponding Lagrange's equations reproduce Lorentz force

$$m \ddot{\vec{r}} = q \left( \vec{E} + \vec{v} \times \vec{B} \right) \quad \begin{array}{l} \text{LORENTZ} \\ \text{FORCE} \end{array}$$

In order to achieve this goal, it is convenient to rewrite the fields in terms of the potentials (scalar and vector)

$$\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t} \quad \vec{B} = \nabla \times \vec{A}$$

One can now prove that the Lagrangian is

$$\begin{aligned} \mathcal{L} &= \frac{1}{2} m \dot{\vec{r}}^2 - qV + q\dot{\vec{r}} \cdot \vec{A} \\ &= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - q \left( V - \dot{x} A_x - \dot{y} A_y - \dot{z} A_z \right) \end{aligned}$$

Indeed one can now check, component by component, that Lagrange's equations reproduce the components of Lorentz equation. Consider for example the  $x$  component

$$\frac{\partial \mathcal{L}}{\partial x} = -q \frac{\partial V}{\partial x} + q\dot{x} \frac{\partial A_x}{\partial x} + q\dot{y} \frac{\partial A_y}{\partial x} + q\dot{z} \frac{\partial A_z}{\partial x}$$

$$\frac{\partial \mathcal{L}}{\partial \dot{x}} = m \dot{x} + q A_x$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}} = m \ddot{x} + q \left( \frac{\partial A_x}{\partial t} + \frac{\partial A_x}{\partial x} \dot{x} + \frac{\partial A_x}{\partial y} \dot{y} + \frac{\partial A_x}{\partial z} \dot{z} \right)$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}} - \frac{\partial \mathcal{L}}{\partial x} = m \ddot{x} + q \left( \frac{\partial A_x}{\partial t} + \cancel{\frac{\partial A_x}{\partial x} \dot{x}} + \frac{\partial A_x}{\partial y} \dot{y} + \frac{\partial A_x}{\partial z} \dot{z} \right) + q \left( \frac{\partial V}{\partial x} - \cancel{\dot{x} \frac{\partial A_x}{\partial x}} - \frac{\partial A_y}{\partial x} \dot{y} - \frac{\partial A_z}{\partial x} \dot{z} \right)$$

$$= m \ddot{x} + q \underbrace{\left( \frac{\partial V}{\partial x} + \frac{\partial A_x}{\partial t} \right)}_{-E_x} + q \dot{y} \underbrace{\left( \frac{\partial A_x}{\partial y} - \frac{\partial A_y}{\partial x} \right)}_{-B_z}$$

$$+ q \dot{z} \underbrace{\left( \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right)}_{B_y} = 0$$

$$m \ddot{x} = q \left( E_x + \underbrace{\dot{y} B_z - \dot{z} B_y}_{\left[ \dot{\mathbf{r}} \times \mathbf{B} \right]_x} \right)$$

The equation above is indeed the x component of Lorentz equation.

It is interesting to evaluate the generalized momentum coming from this Lagrangian

$$p_x = \frac{\partial \mathcal{L}}{\partial \dot{x}} = m \dot{x} + q A_x \rightarrow \bar{p} = m \bar{v} + q \bar{A}$$

This is useful in quantum mechanics because the quantization rules require to replace the momentum with the operator

$$\bar{p} \xrightarrow{\text{Q.M.}} -i \hbar \nabla$$

$$\text{Therefore } m \bar{v} \xrightarrow{\text{Q.M.}} -i \hbar \nabla - q \bar{A}$$